

average surface radar properties best reflect the character of the various venusian plains. The mean planetary radius is 6051.9 km, a value slightly above that previously reported [Masursky *et al.*, 1980; Pettengill *et al.*, 1982]; we have referenced all elevations to a datum of 6051.0 km in order to be compatible with topographic provinces defined by Masursky *et al.* [1980]. Table 1 compares the topographic provinces proposed by Masursky *et al.* [1980] with  $\rho$  values for each elevation interval.

The  $\alpha$  values are derived from fitting discrete templates to the altimetric measurements in a two-stage process [Pettengill *et al.*, 1980a; Garvin *et al.*, 1984c]. These templates are approximations to the observed radar echo distribution as a function of time delay. Templates for a wide variety of altitudes and roughnesses ( $\alpha$ ) were used in a weighted least squares comparison to obtain good estimates of the radar parameters [Pettengill *et al.*, 1980a]. Unlike roughness,  $\rho$  is linear in equation (1) and can be fit in a continuous sense to the observed data. Surfaces poorly fit by the templates because of their nonstandard quasi-specular scattering behavior often have  $\rho$  values with low relative errors. This is because  $\rho$  serves only as a scaling parameter and does not affect the shape of the quasi-specular scattering curve. As such,  $\rho$  is a relatively model-independent parameter, in contrast with  $\alpha$ , which totally defines the shape of the curve, and is thus model dependent.

#### GEOLOGIC INTERPRETATION

In order to interpret the radar  $\rho$  in terms of the material properties of a surface, several assumptions must be made. We adopt a single-layer model with a magnetic permeability of unity in order that  $\rho$  is a function of the complex dielectric discontinuity at the reflecting interface (typically the uppermost surface) [Pettengill *et al.*, 1982]. It is then possible to express  $\rho$  in terms of the complex dielectric constant  $\epsilon$  as follows

$$\rho = |(1 - \epsilon^{1/2}) / (1 + \epsilon^{1/2})|^2 \quad (2)$$

where

$$\epsilon = \kappa + j\{L + S/(\epsilon_0\omega)\} \quad (3)$$

for  $j = (-1)^{1/2}$ . The real component  $\kappa$  dominates the radar  $\rho$  for dry lunar and terrestrial rocks and soils [Krotikov and Troitsky, 1963; Campbell and Ulrichs, 1969; Olhoeft and Strangway, 1975; Pettengill *et al.*, 1982]. The imaginary component  $\{L + S/(\epsilon_0\omega)\}$  depends on the dielectric loss  $L$ , the volume conductivity  $S$ , the permittivity of free space  $\epsilon_0$ , and the angular frequency of the incident radiation  $\omega$  (here  $\omega$  is 1.76 GHz); for our purposes none of these factors affect  $\epsilon$  very much, and  $\epsilon \approx \kappa$  for typical terrestrial and lunar geologic materials [Olhoeft and Strangway, 1975; Pettengill *et al.*, 1982].

The bulk dielectric constant  $\kappa$  is directly dependent upon the porosity and electrical properties of geologic materials, and indirectly on composition (e.g., Fe and Ti content) and bulk density  $\gamma$ . The dielectric properties of typical terrestrial and lunar rocks, soils, and minerals at radar frequencies ( $\sim 10^9$  Hz or 3.8–70 cm wavelengths) have been determined [Krotikov, 1962; Krotikov and Troitsky, 1963; Evans and Hagfors, 1968; Campbell and Ulrichs, 1969; Olhoeft and Strangway, 1975] and empirical relationships between these properties and density [Krotikov and Troitsky, 1963; Olhoeft and Strangway, 1975] as well as composition [Hansen *et al.*, 1973; Saint-Amant and Strangway, 1970] have been derived. Krotikov

[1962] and Krotikov and Troitsky [1963] found that for a wide range of naturally occurring rocks and soils (e.g., from tuff to dunite), the relationship

$$\gamma = (\kappa^{1/2} - 1)/0.5 \quad (4)$$

for density  $\gamma$  and dielectric constant  $\kappa$  was accurate to within 10% at 4 cm and 12 cm radar wavelengths. All samples were dried. Equation (4) can be related to radar  $\rho$  via equation (2) so that

$$\gamma = 4\rho^{1/2}/(1 - \rho^{1/2}) \quad (5)$$

or, inversely,

$$\rho = \gamma^2 / \{4(2 + 0.5\gamma)^2\} \quad (6)$$

Note that this is a quadratic relationship that empirically applies only for dry rocks and soils with  $\gamma$  between 1.5 and 3.3 g/cm<sup>3</sup>. Campbell and Ulrichs [1969] measured  $\kappa$  as a function of  $\gamma$ ,  $T$  (temperature), porosity (i.e., powder versus solid rock), and frequency (0.45 and 35 GHz). They observed that the Rayleigh mixing formula well modeled the change in  $\kappa$  when the density of a solid was reduced by powdering the sample. Though they did not derive an equation relating bulk density  $\gamma$  and  $\kappa$ , their data fit the Krotikov [1962] model (equation (4)) very well (correlation  $R^2 > 0.90$ ).

Olhoeft and Strangway [1975] analyzed a large sample of lunar rocks and soils in terms of  $\kappa$  and derived an empirical relationship between  $\gamma$  and  $\kappa$  for frequencies on the order of  $10^6$  Hz. For dried samples they found the following power law relationship:

$$\kappa = a^\gamma \quad (7)$$

where  $a = 1.93 \mp 0.17$  and  $\gamma$  is density in grams per cubic centimeter. This can be written in terms of radar  $\rho$  as follows

$$\gamma = (2/\ln a) \ln \{(1 + \rho^{1/2}) / (1 - \rho^{1/2})\} \quad (8)$$

In order to evaluate differences between the relationships described by equations (5) and (8) and to choose the relationship most appropriate for Venus, a statistical correlation analysis is required. Correlation of the Krotikov [1962] relationship (equation (5)) with that of Olhoeft and Strangway [1975] using a least squares regression technique demonstrates that

$$\rho_k = 0.64\rho_{os} + 0.54 \quad (9)$$

at a correlation  $R^2$  of 0.99 with  $a = 1.87$ , where  $\rho_k$  is from equation (5) and  $\rho_{os}$  is from equation (8). If only solid rocks are considered, the correlation can be improved. In equation (9) the radar  $\rho$  values are equivalent for  $\gamma = 2.8$  g/cm<sup>3</sup> at  $\rho = 0.17$ . On the basis of the above analysis, we adopt the relationships

$$\kappa = 1.87^\gamma \quad (10)$$

$$\gamma = 3.2 \ln \{(1 + \rho^{1/2}) / (1 - \rho^{1/2})\} \quad (11)$$

as those to use when interpreting the radar  $\rho$  in terms of density  $\gamma$ . We have chosen to use this combined relationship in favor of either of the two independent models (equations (5) and (8)) because it links the best aspects of both relationships in a single, empirical equation (equation (11)) for the range of  $\gamma$  values of geologic interest.

From equation (11) the following simplified interpretation of the magnitude of reflectivity  $\rho$  can be derived: (1)  $0.02 < \rho < 0.10$  correspond to  $\gamma < 2.0$  g/cm<sup>3</sup> and are likely to be high-porosity materials such as soil, poorly welded tuff, or weakly cemented sediments, (2)  $0.10 < \rho < 0.20$  correspond to