



Fig. 14. Observed cratering efficiency corrected for ambient pressure  $\pi_p$  and acceleration  $g$  ( $d/g$ ) as given by equation (18) and expressed as functions of variables controlling dynamic pressure  $P_w$  (Figure 14a) and the stagnation pressure  $P_s$  of the impactor (Figures 14b and 14c) as inferred from Figure 13. Figure 14a includes data for a limited range of Reynolds numbers (1-14) of the ejecta where  $C_D \sim 1/Re$  and low Mach numbers ( $M < 6$ ) for no. 140-200 sand and compacted pumice. Both targets exhibit enhancement of cratering efficiency above a critical value of the subsonic wake dynamic pressure ( $P_w = \rho v_i^2$ ). Figure 14b considers hypervelocity at high Mach numbers ( $M > 6$ ) and the assumption that the wake collides at supersonic velocities. Under such conditions, the wake resembles an impinging compressible gas jet, and the stagnation pressure  $P_s$  becomes  $PM^2 \gamma (\gamma - 1)$ . Both sand and pumice exhibit relative enhancement of cratering efficiency over conditions minimizing wake effects. If the wake back pressure completely offsets ambient pressure effects, then the dimensionless wake

pressure parameter might be expected to follow the same exponent but with opposite sign, i.e.,  $\beta' \alpha \sim 0.3$  as shown. The data shown are for a limited range of Reynolds numbers: 1-6 for pumice, 6-14 for sand. Figure 14c considers just pumice data over two restricted ranges of Reynolds numbers, thereby resulting in  $d/g$  becoming proportional to  $C_D \rho R_w / \delta_t \rho$ , where  $C_D$  is a constant (here  $d/g$  is taken simply as  $\rho R_w$ , with  $\rho$  referenced to air at one bar). At lower Reynolds numbers ( $\sim 6$ ), cratering efficiency appears to be enhanced with increasing stagnation pressure consistent with Figure 14b. But at higher Reynolds numbers (10), it decreases. This may reflect a change from wake effects decoupled from the impactor to wake effects coupled to the impactor, thereby modifying the effective energy density. (Abscissas are shown for density and pressure referenced to air at 1 bar with velocities in km/s and target density in cgs; the ordinate is in dimensionless form. Figure 14c uses  $C_D = 1$ .)